

# **Using Higher Moments of Fluctuations and their Ratios in the Search for the QCD Critical Point**

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**work with: Krishna Rajagopal (MIT)**

**Misha Stephanov (University of Illinois)**



Massachusetts Institute of Technology

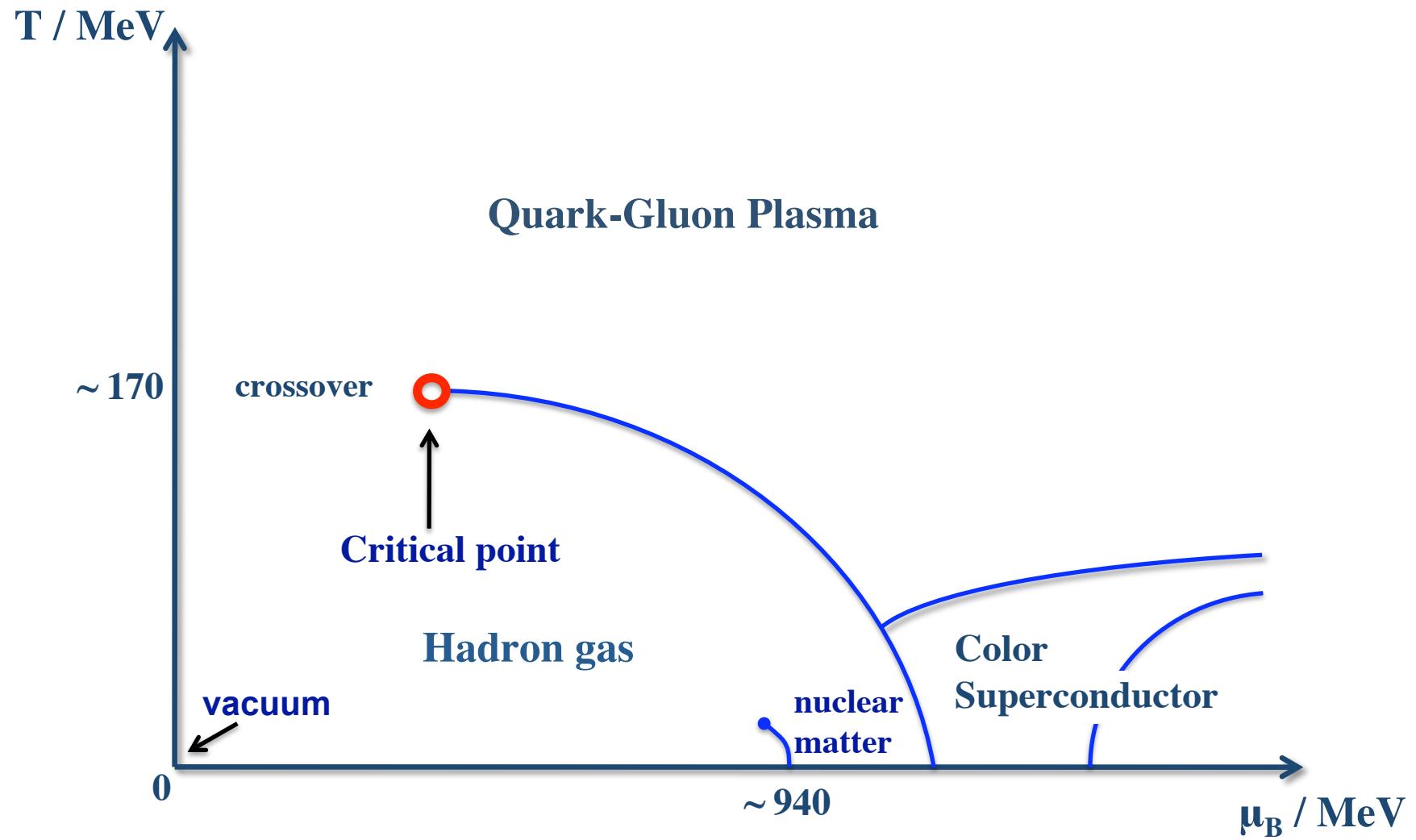
# Outline

- Introduction
- Critical Contribution to Particle Multiplicity Fluctuations
- Ratios of Fluctuation Observables
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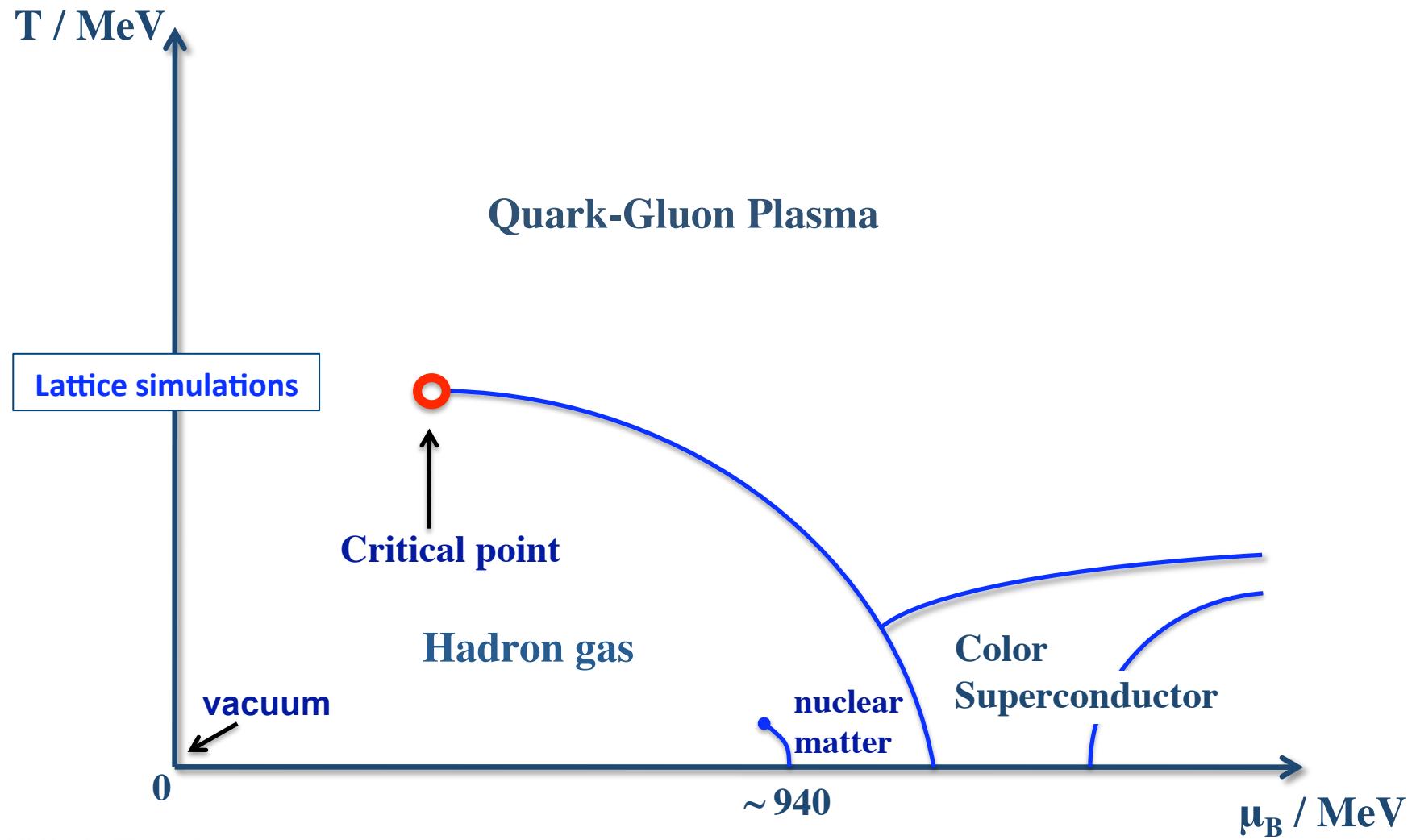
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# QCD Phase Diagram



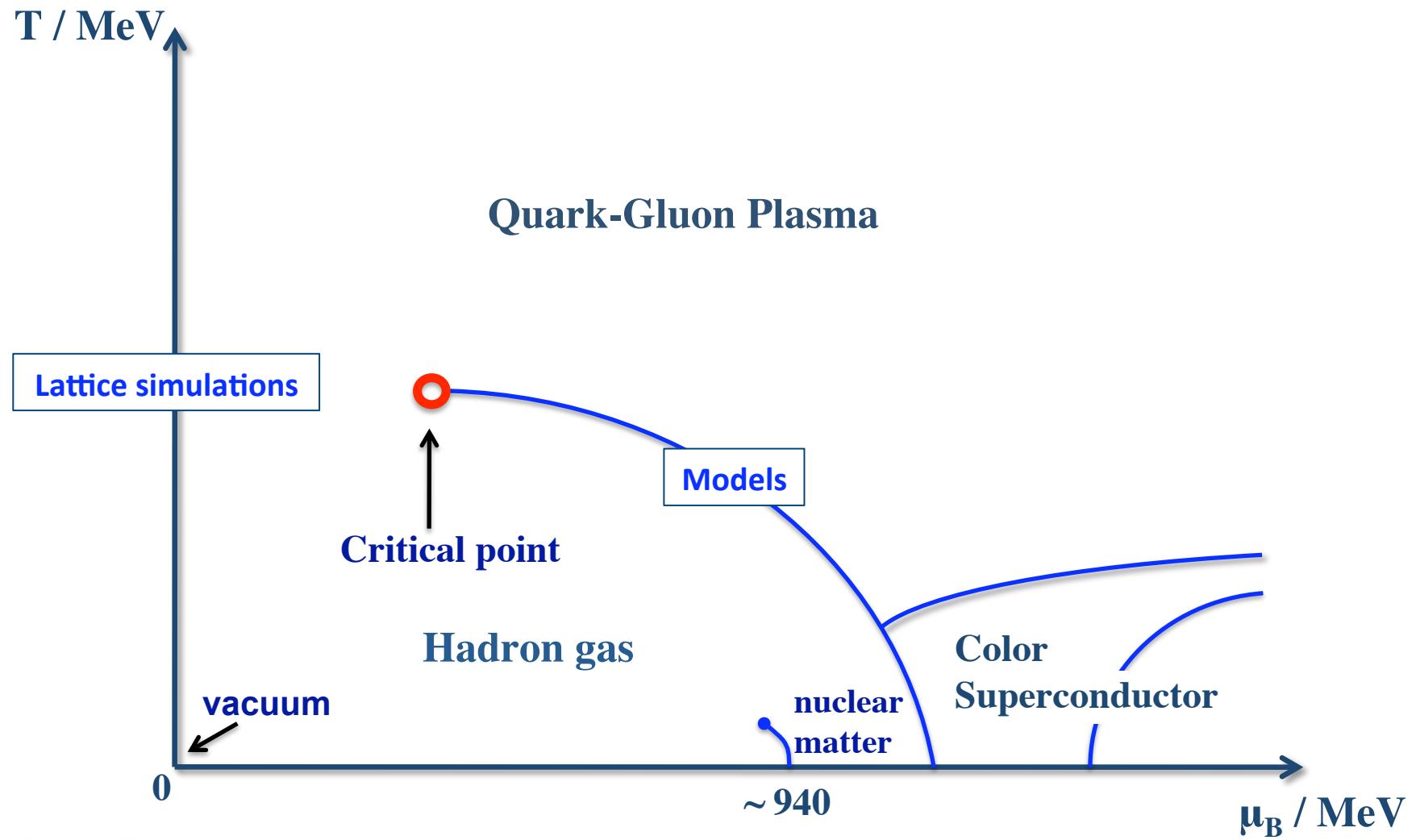
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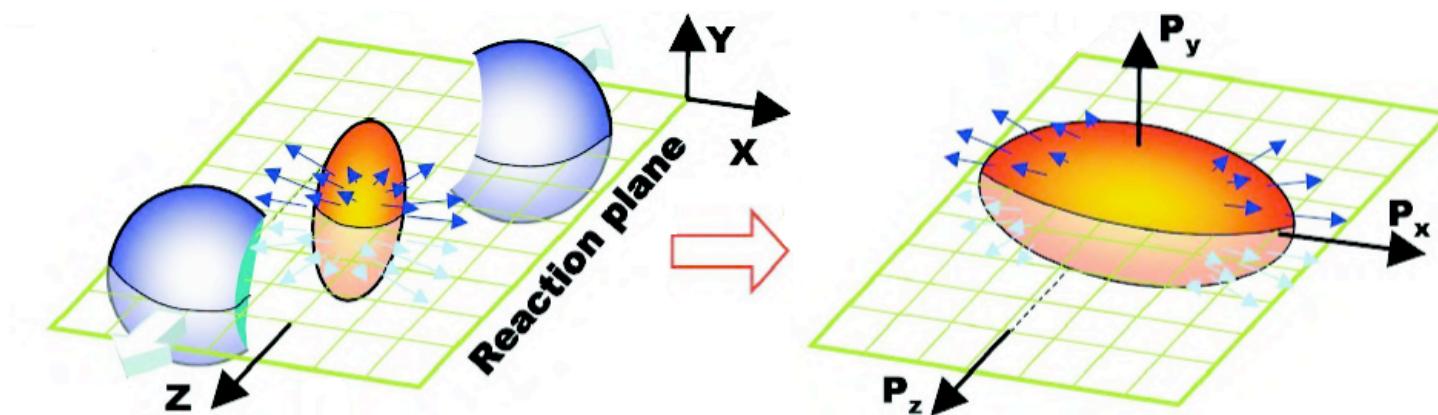
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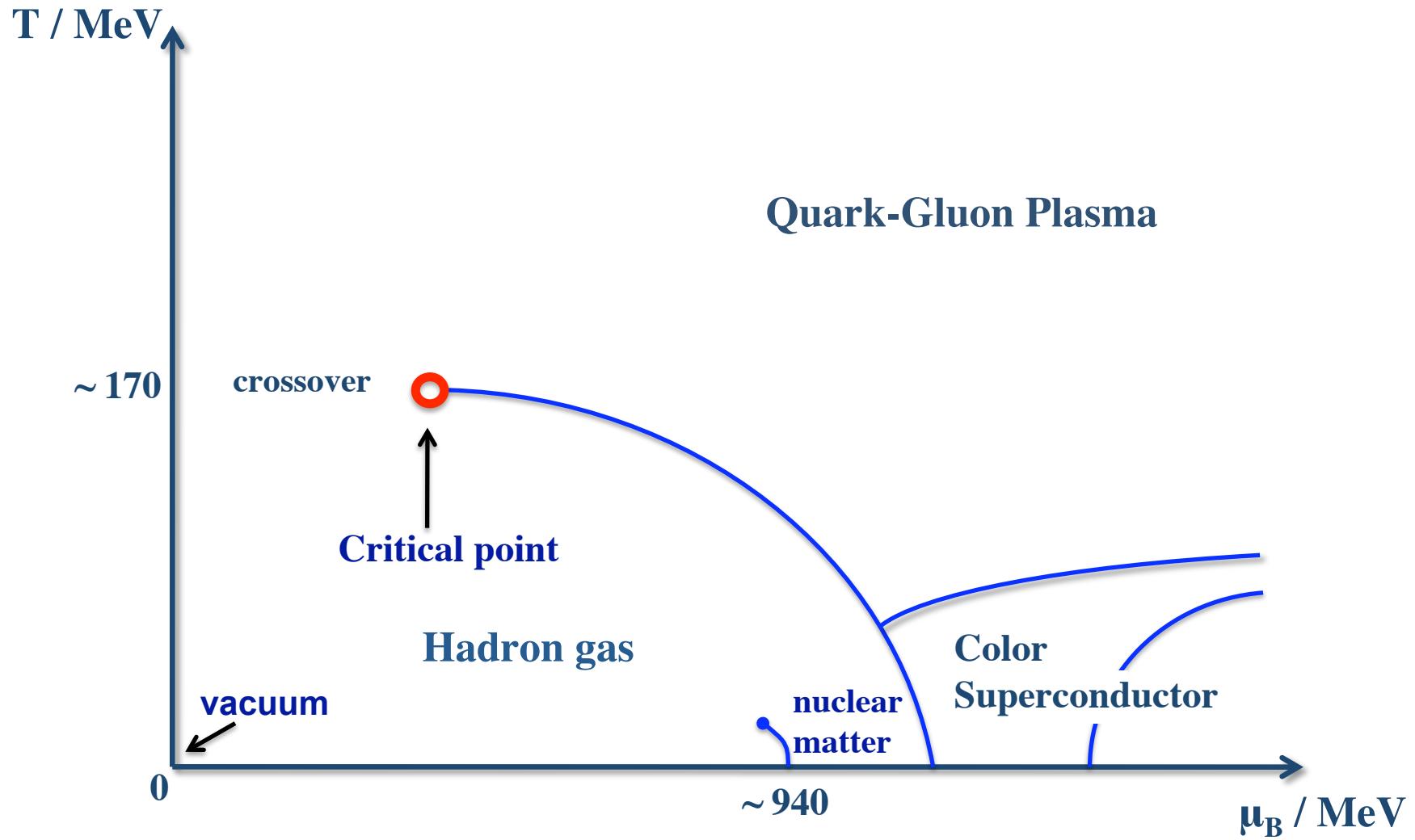
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- Momentum asymmetry → collective flow  
→ strongly-coupled QGP

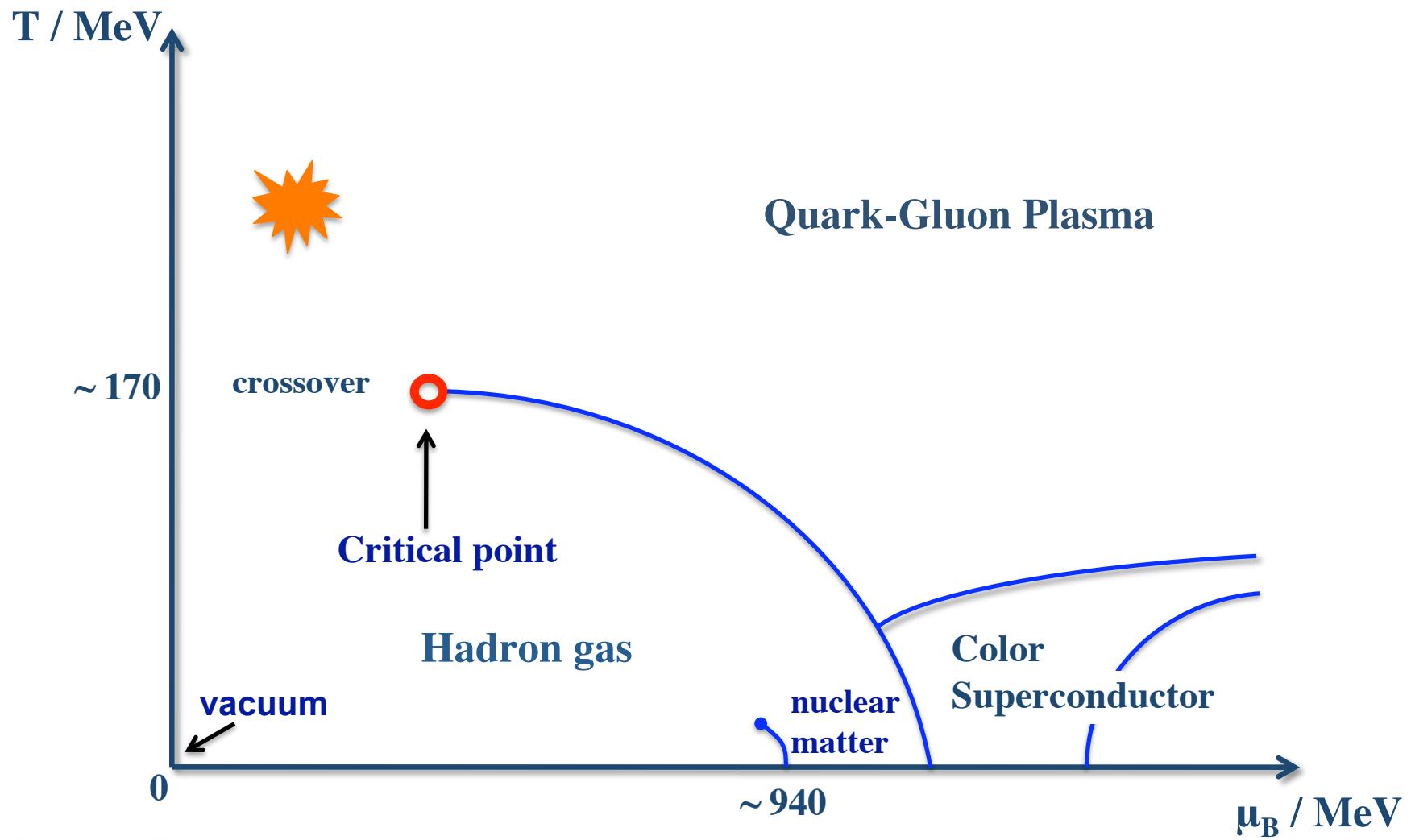


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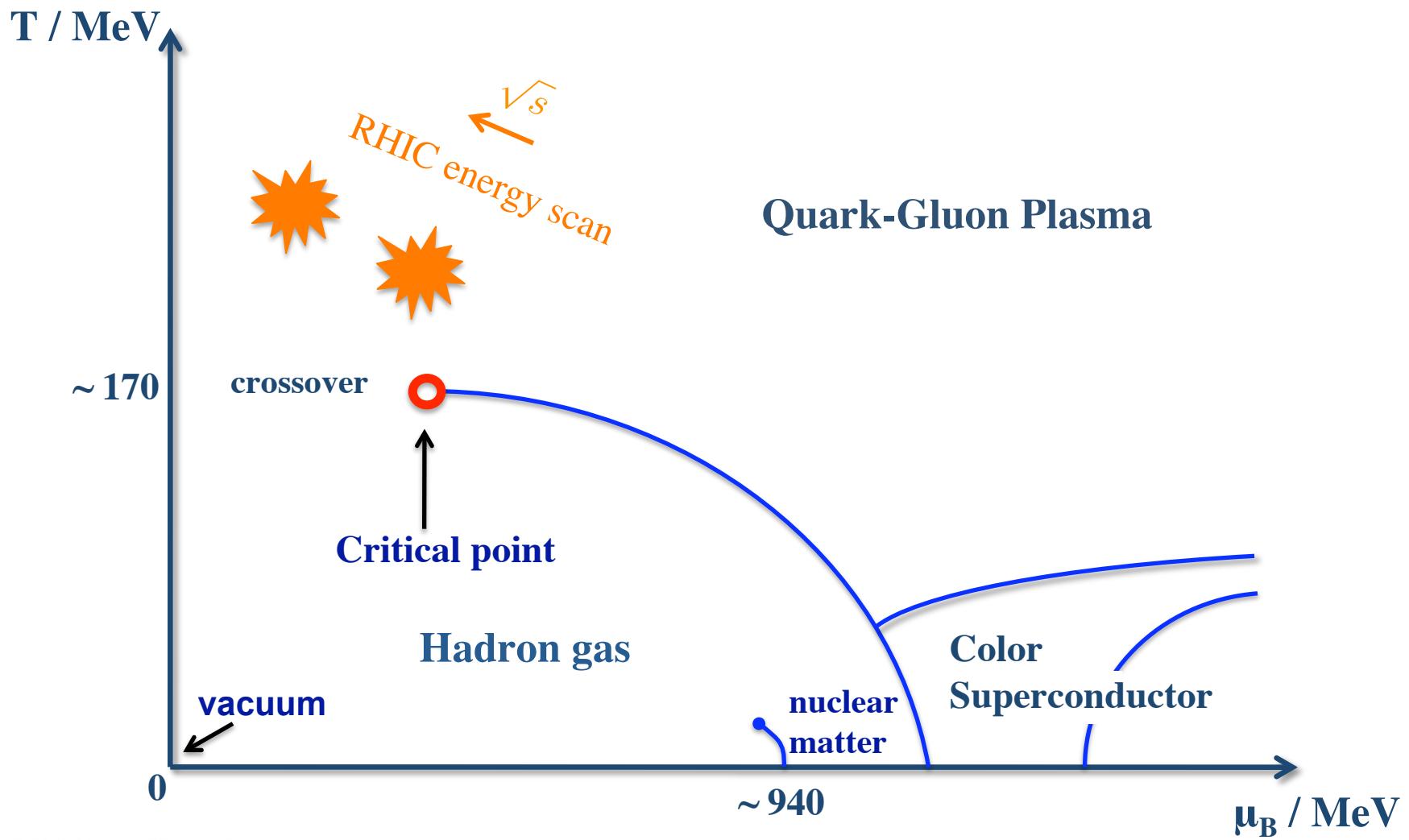
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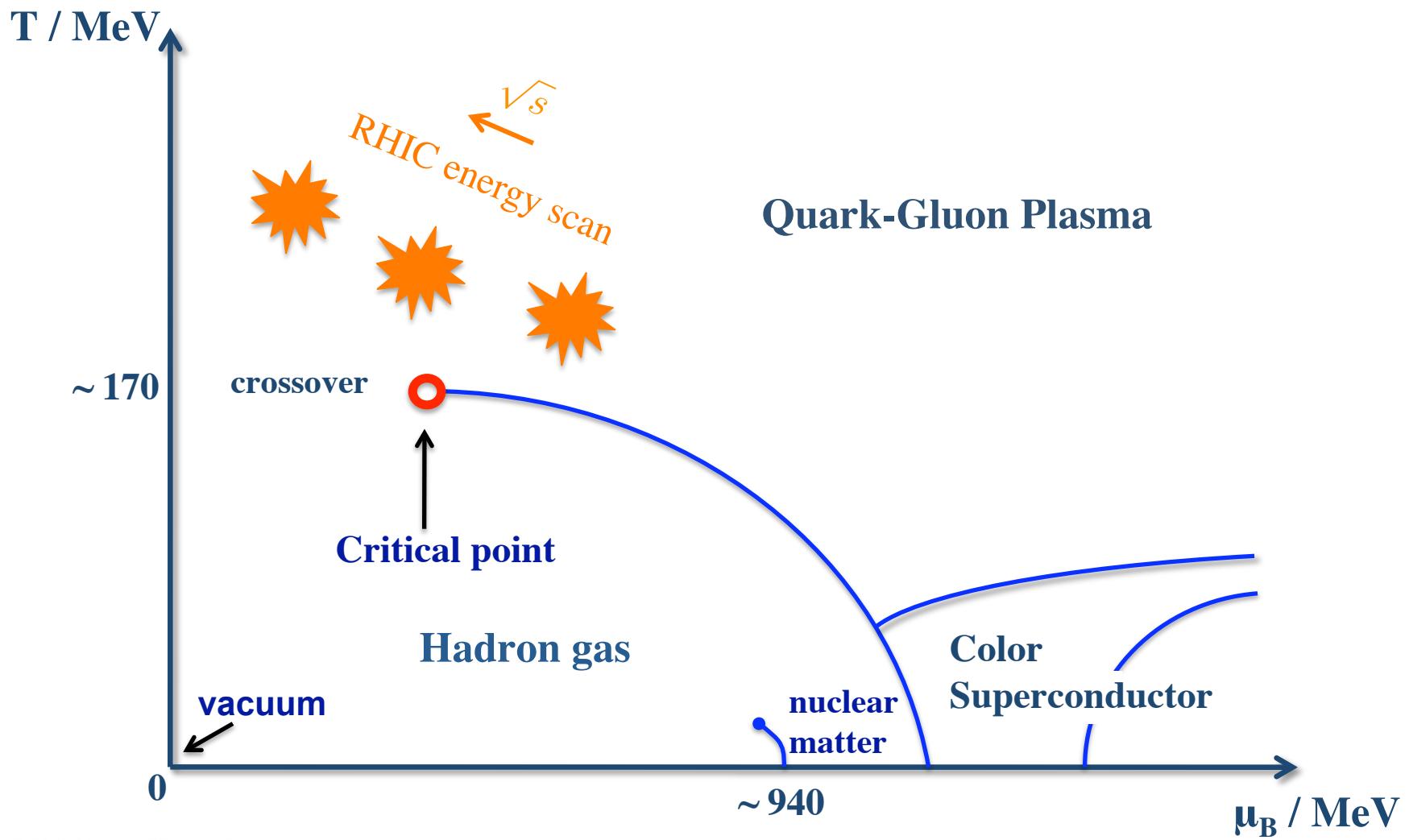


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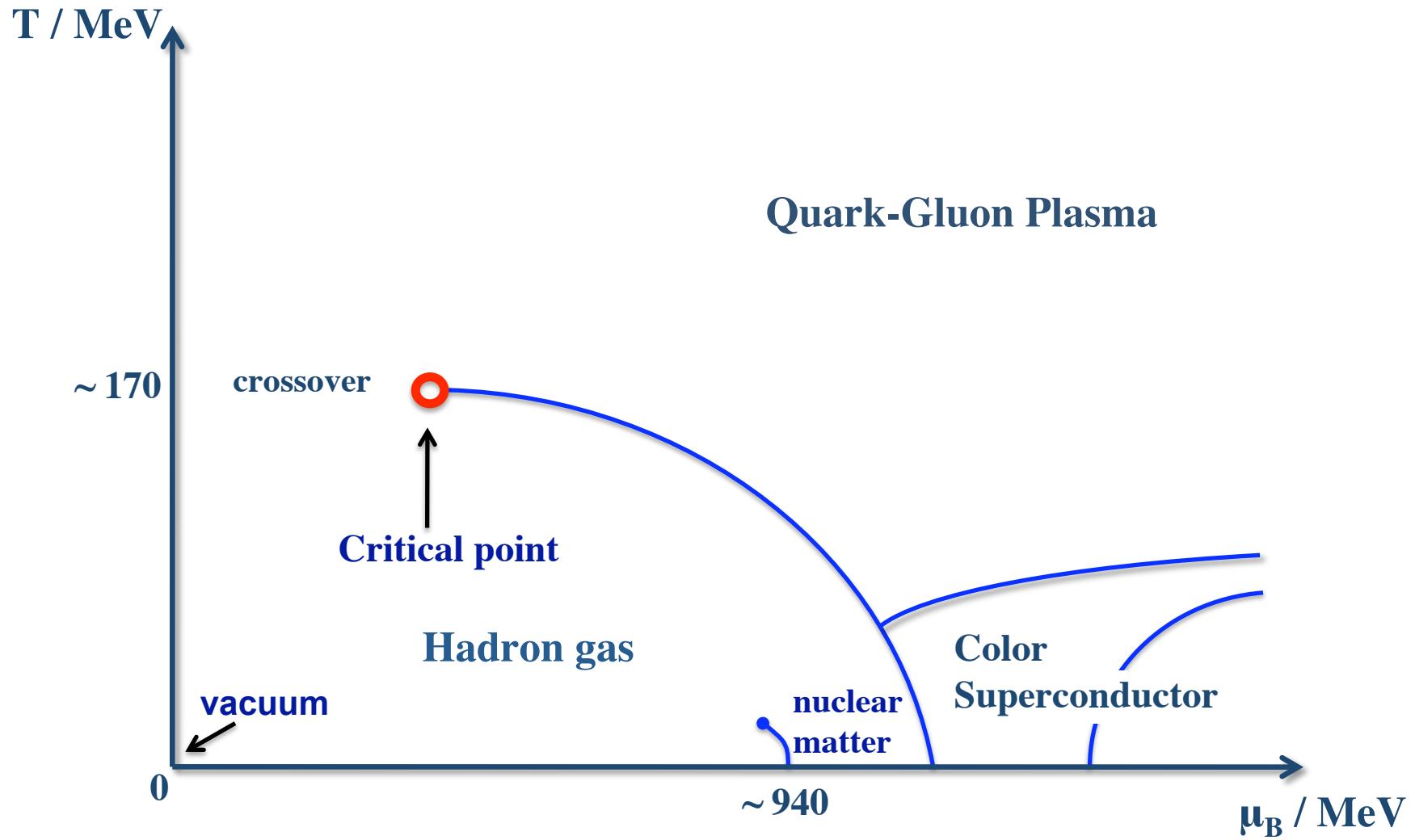
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# Heavy-Ion Collision Experiments - continued

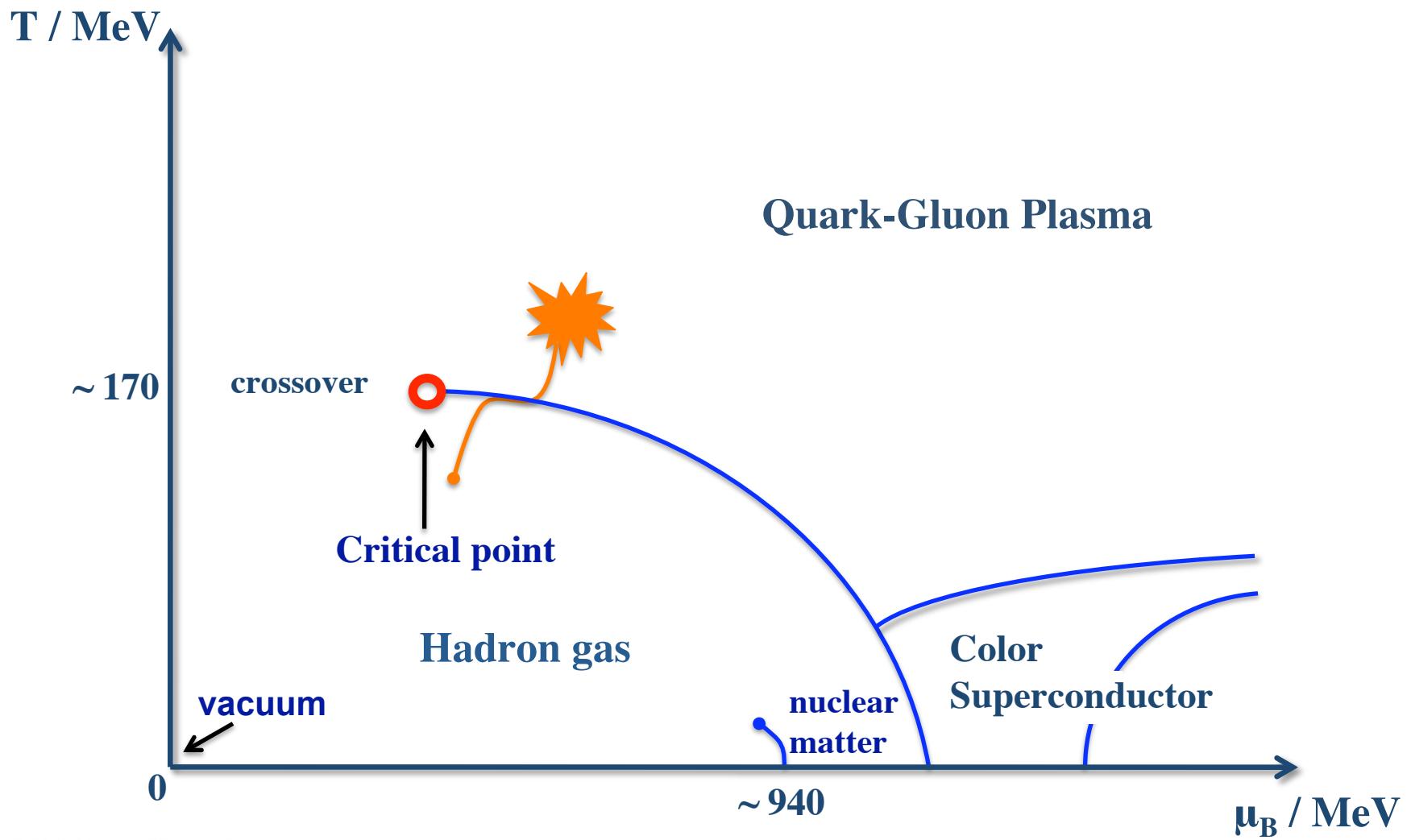
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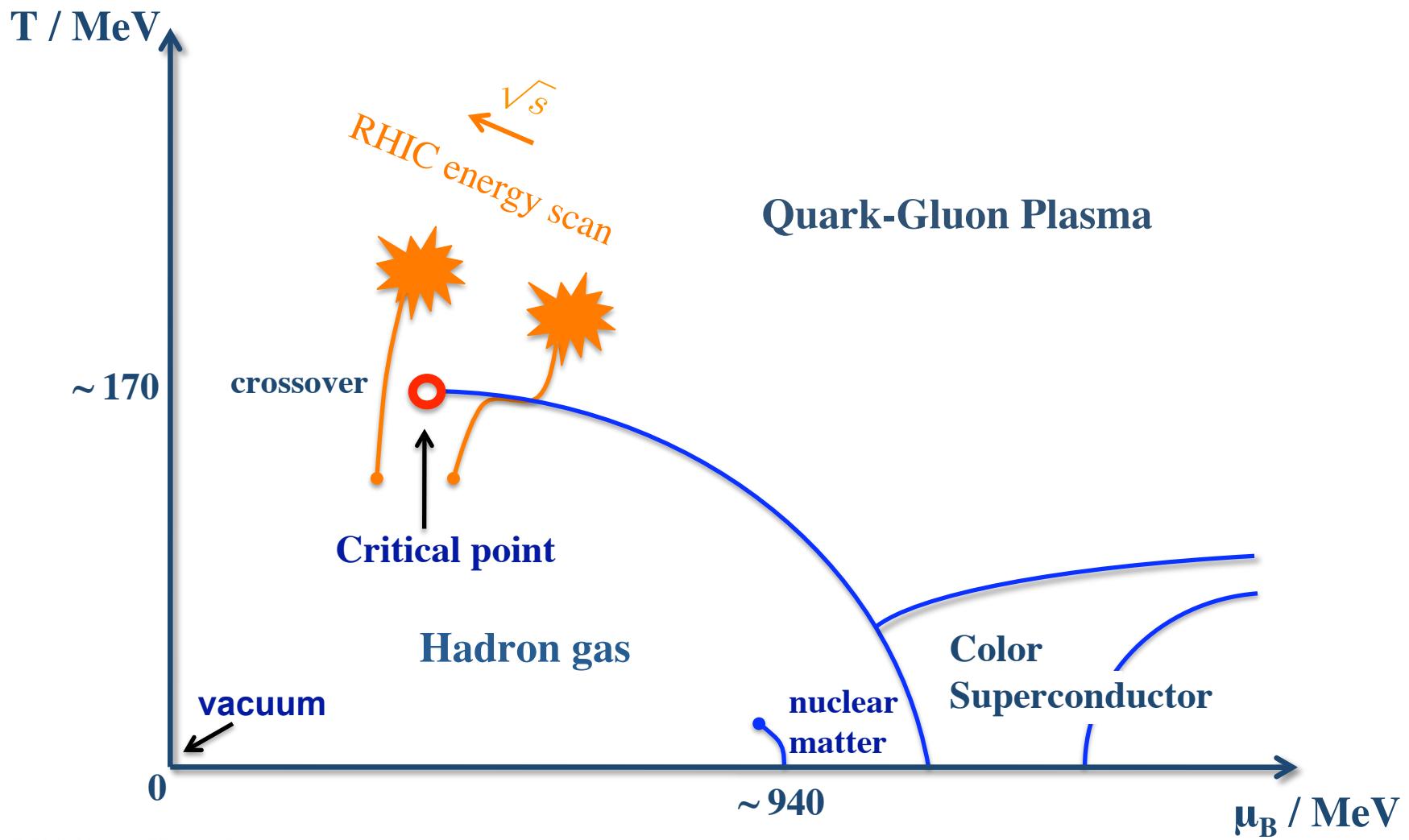
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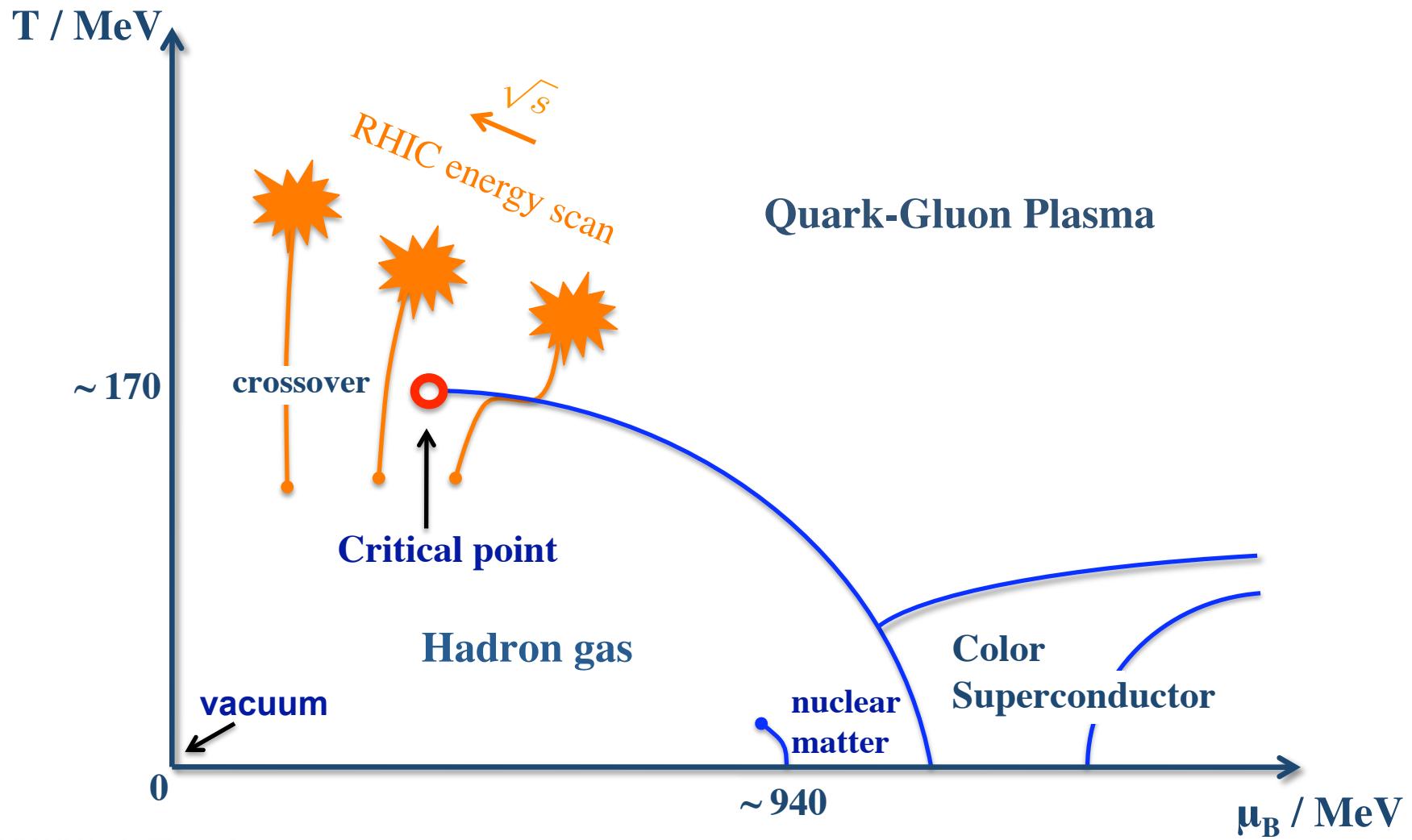
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- Event-by-Event fluctuations

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- Correlation length  $\xi = m_\sigma^{-1}$  diverges at the CP
- Near the CP:  $\lambda_3 = \tilde{\lambda}_3 T (T \xi)^{-3/2}$ ,  $\lambda_4 = \tilde{\lambda}_4 (T \xi)^{-1}$  with  $0 \lesssim \tilde{\lambda}_3 \lesssim 8$ ,  $4 \lesssim \tilde{\lambda}_4 \lesssim 20$  dimensionless and known in the Ising universality class

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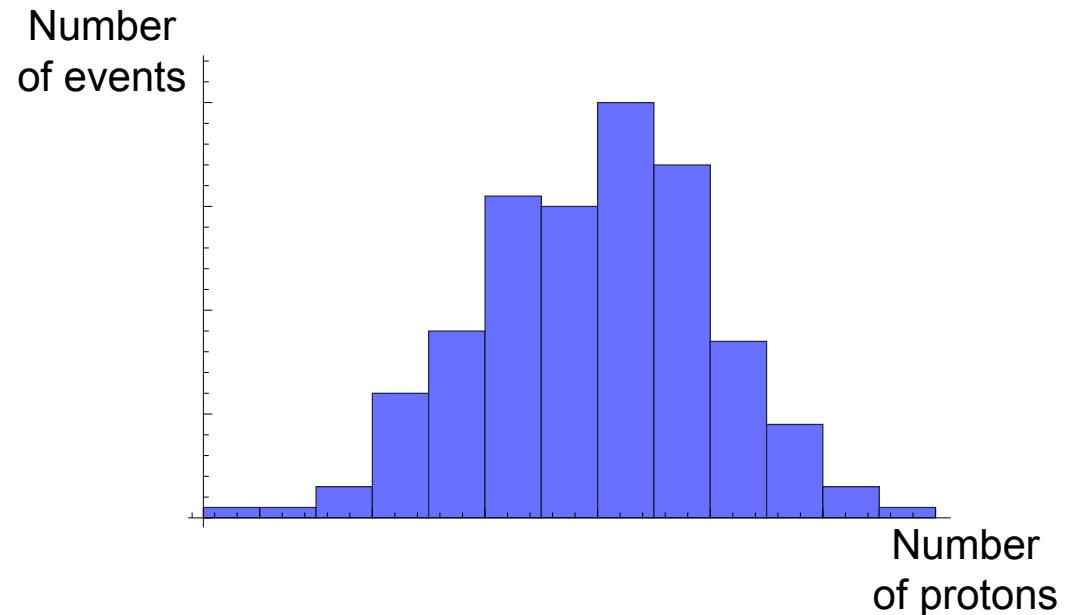
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  - Particle multiplicity fluctuations
  - Momentum distributions
  - Ratios, etc...of these particles.

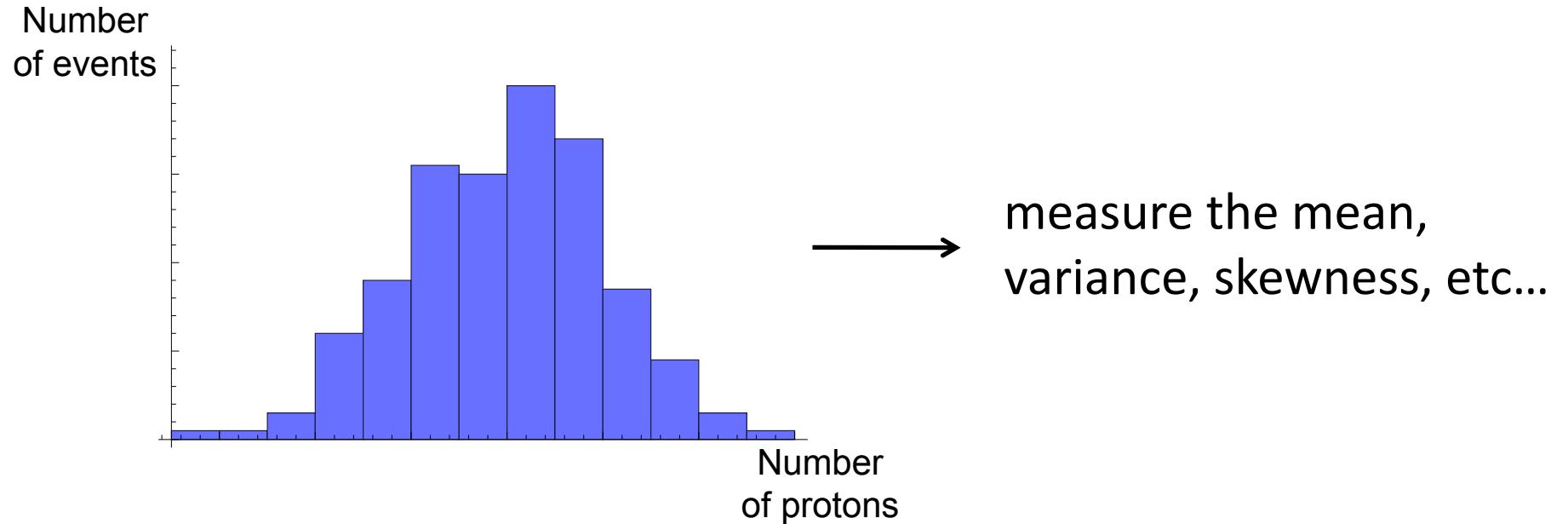
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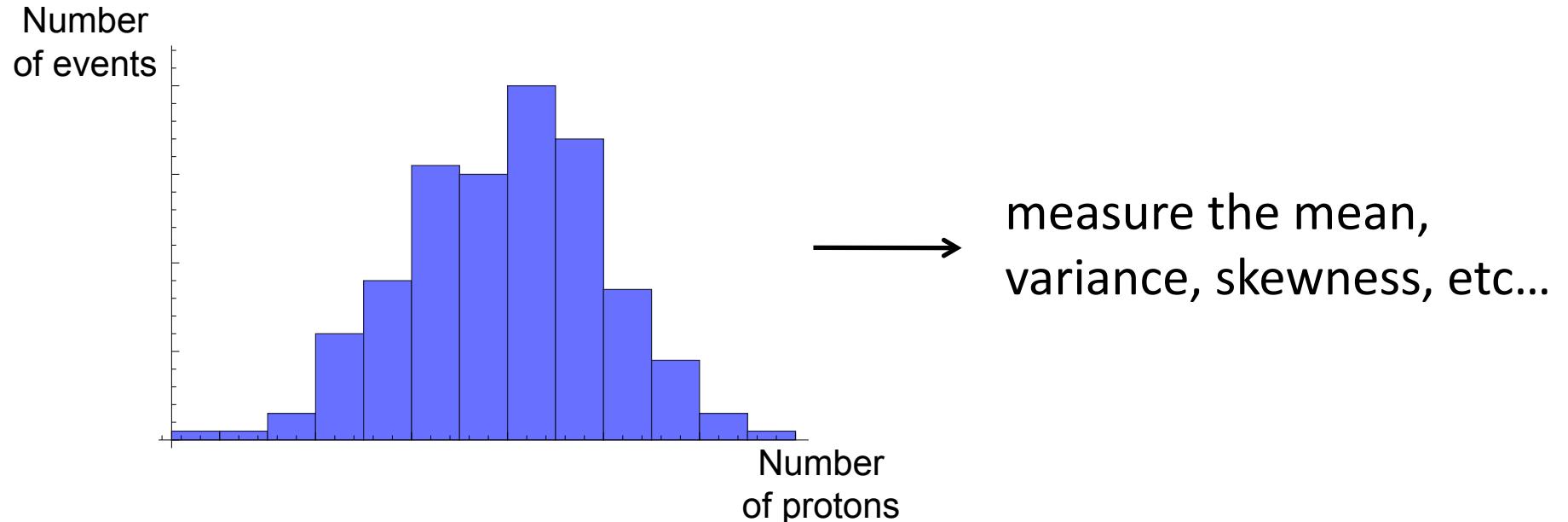
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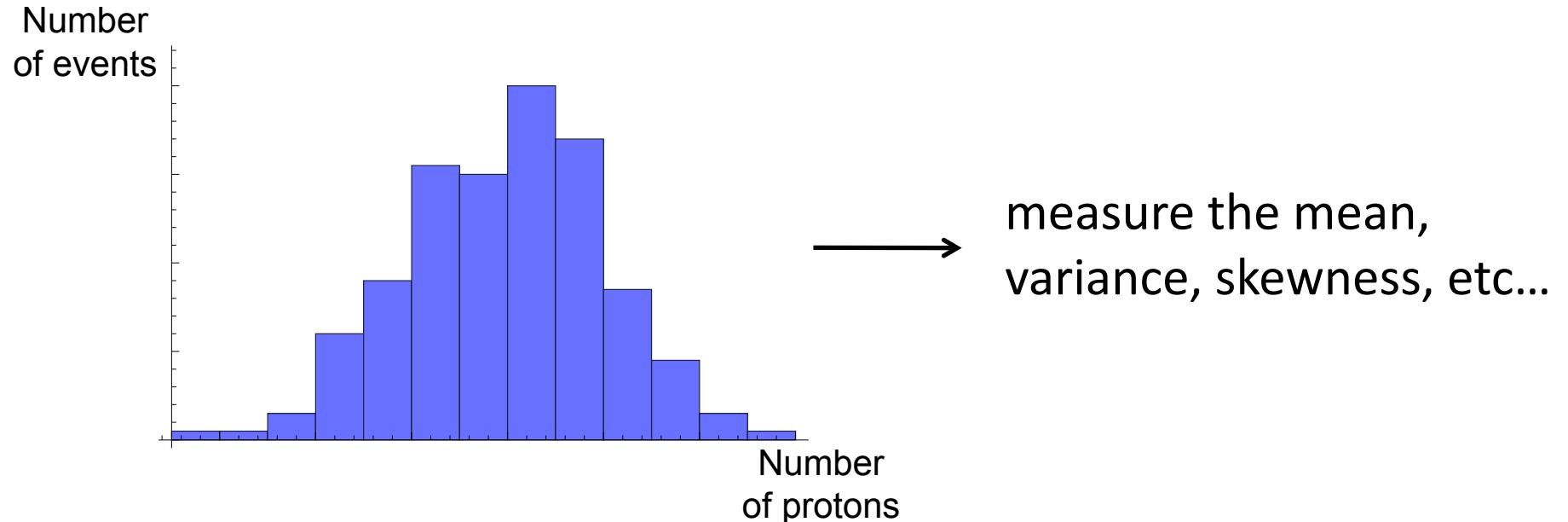


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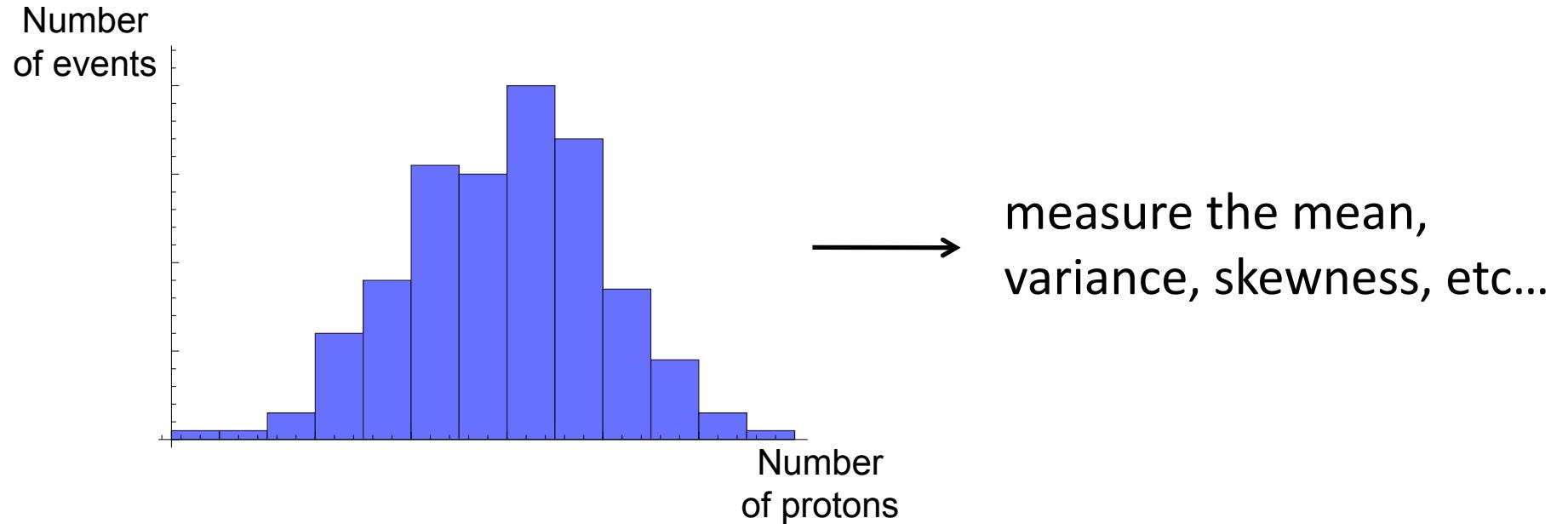
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- Want to obtain the critical contribution to these quantities
- We will use cumulants, e.g.:

$$\kappa_2 = \langle N^2 \rangle, \quad \kappa_3 = \langle N^3 \rangle, \quad \kappa_4 \equiv \langle\langle N^4 \rangle\rangle = \langle N^4 \rangle - 3\langle N^2 \rangle^2$$

# Critical contribution to pion/proton correlators



$$\langle \delta n_{\mathbf{k}_1} \delta n_{\mathbf{k}_2} \rangle_\sigma = d^2 \frac{1}{m_\sigma^2 V} \frac{g^2}{T} \frac{v_{\mathbf{k}_1}^2}{\gamma_{\mathbf{k}_1}} \frac{v_{\mathbf{k}_2}^2}{\gamma_{\mathbf{k}_2}}$$

$$m_\sigma = \xi^{-1}, \quad \gamma_{\mathbf{k}} = \sqrt{\mathbf{k}^2 + m^2}/m, \quad v_{\mathbf{k}}^2 = n_{\mathbf{k}}(1 \pm n_{\mathbf{k}}),$$

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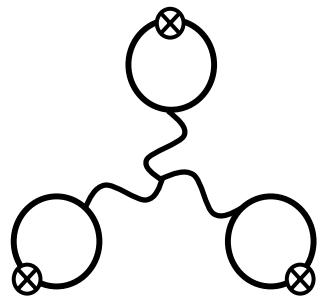
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(Rajagopal, Shuryak, Stephanov 99, Stephanov 08)

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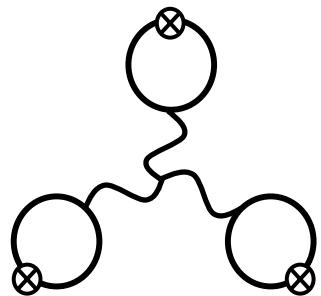
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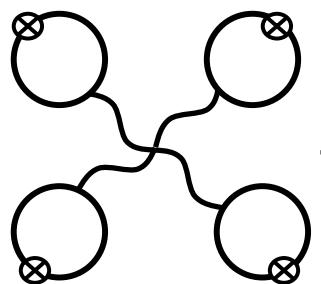
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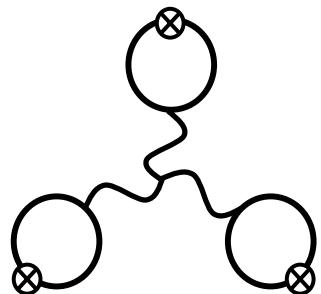
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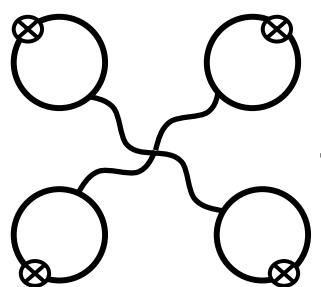
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- Note: correlators depend on 5 parameters:

$$G, g, \xi, \tilde{\lambda}_3, \tilde{\lambda}_4$$

which have large uncertainties

# Calculating multiplicity cumulants

- Second cumulant – variance:

$$\kappa_{2p,\sigma} = \langle (\delta N_p)^2 \rangle_\sigma = \int_{\mathbf{k}_1} \int_{\mathbf{k}_2} \langle \delta n_{\mathbf{k}_1} \delta n_{\mathbf{k}_2} \rangle_\sigma \propto V^1$$

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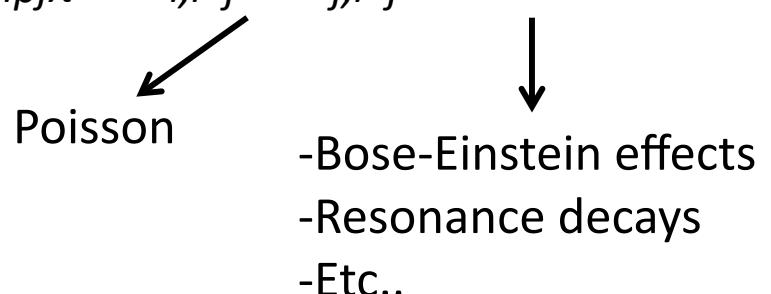
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Poisson

ignore

-Bose-Einstein effects  
-Resonance decays  
-Etc..

# Multiplicity cumulants – critical point signature

- Higher cumulants depend stronger on  $\xi$ :  
 $\omega_2 \propto \xi^2$ ,  
 $\omega_3 \propto \xi^{9/2}$ ,  
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 $\omega_3 \propto \xi^{9/2}$ ,  
 $\omega_4 \propto \xi^7$
- As we approach the CP  $\xi$  increases and then decreases as we move away from it

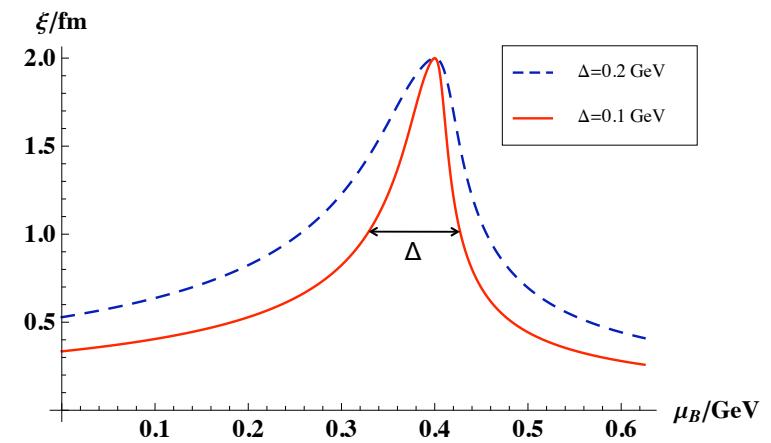
# Multiplicity cumulants – critical point signature

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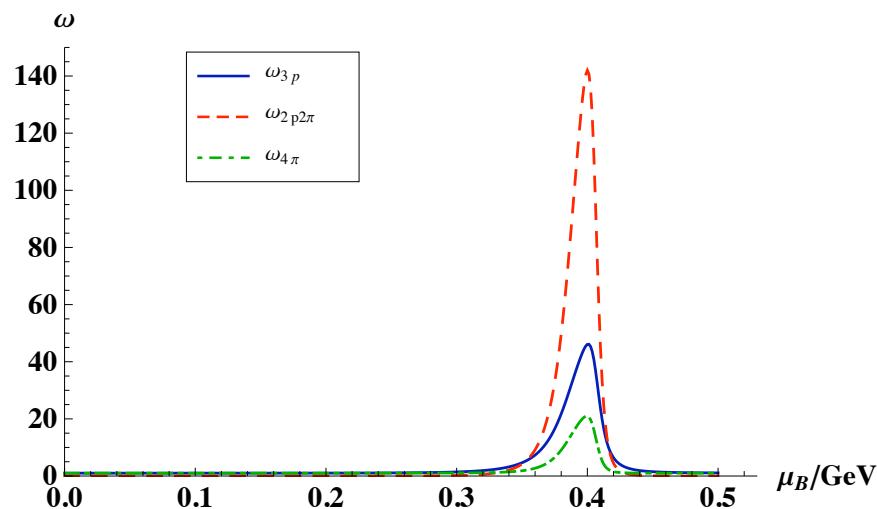
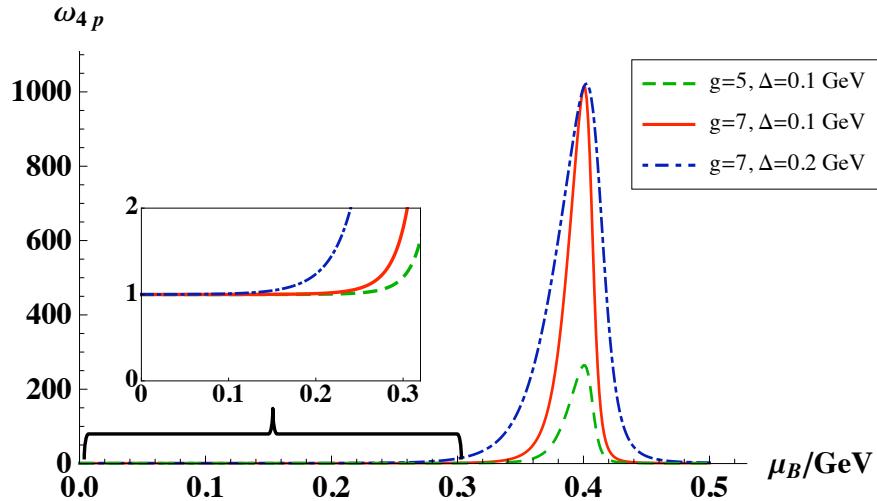
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- **CP signature: Non-monotonic behavior, as a function of collision energy, of multiplicity cumulants**
- E.g. toy example

$$\xi(\mu_B) = \frac{2 \text{ fm}}{(1 + (\mu_B - 400)^2/W^2)^{1/3}}$$

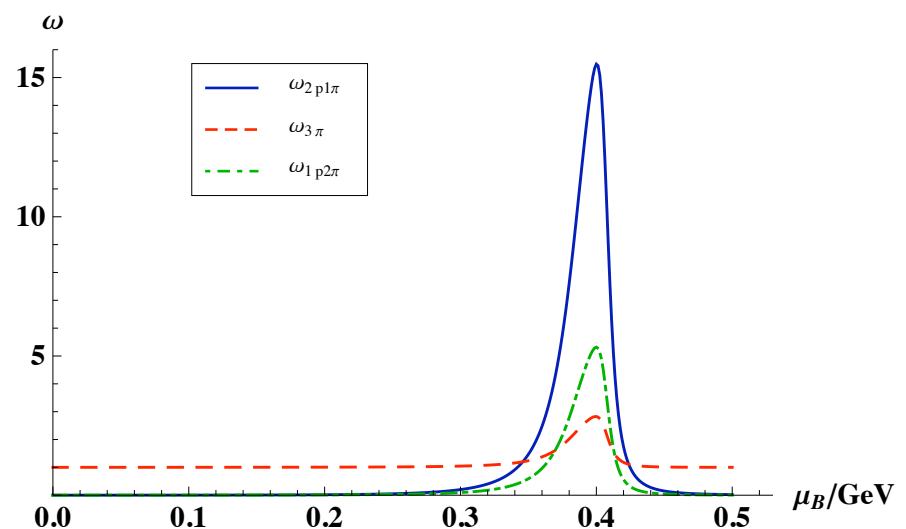


# Multiplicity cumulants – example plots

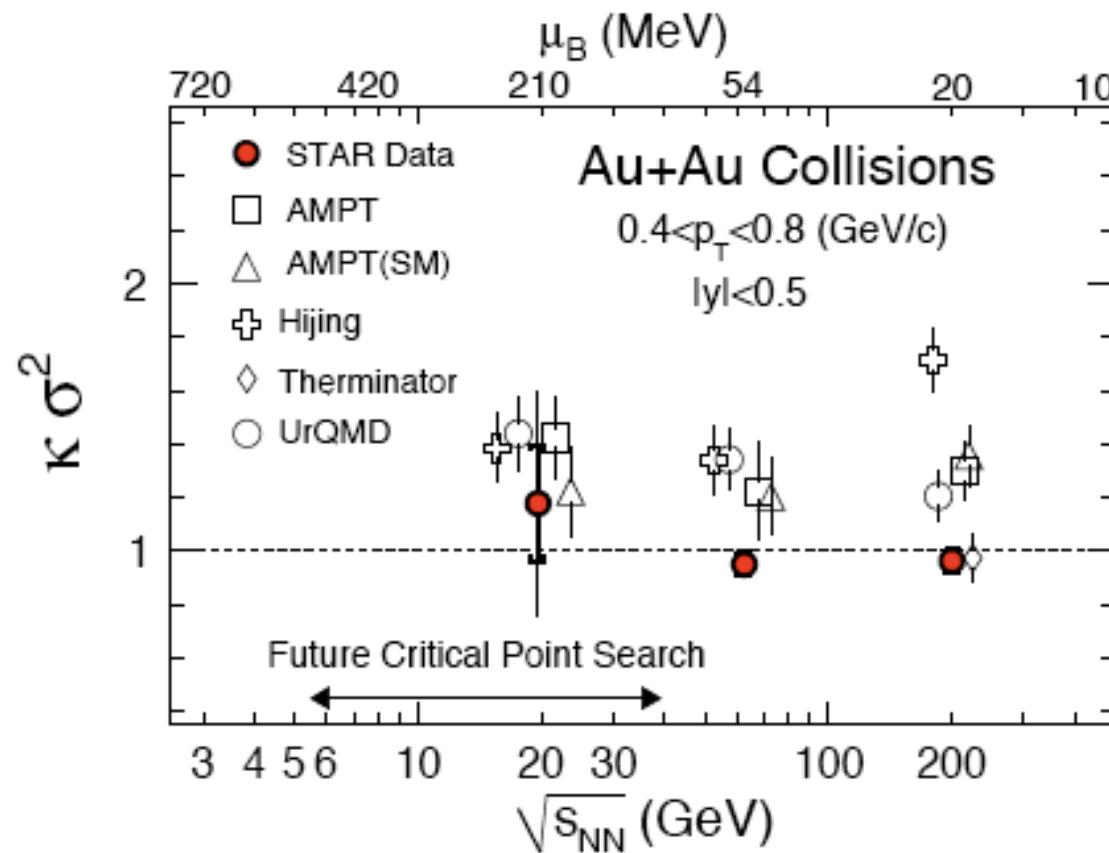


Parametrization (Cleymans et al 05):  

$$T(\mu_B) = a - b\mu_B^2 - c\mu_B^4$$
  
 $a = 0.166 \text{ GeV}, \quad b = 0.139 \text{ GeV}^{-1}, \quad c = 0.053 \text{ GeV}^{-3}$   
and using  
 $\tilde{\lambda}_3 = 6, \quad \tilde{\lambda}_4 = 22, \quad G = 300 \text{ MeV}, \quad g = 7$



# Data on net proton cumulants



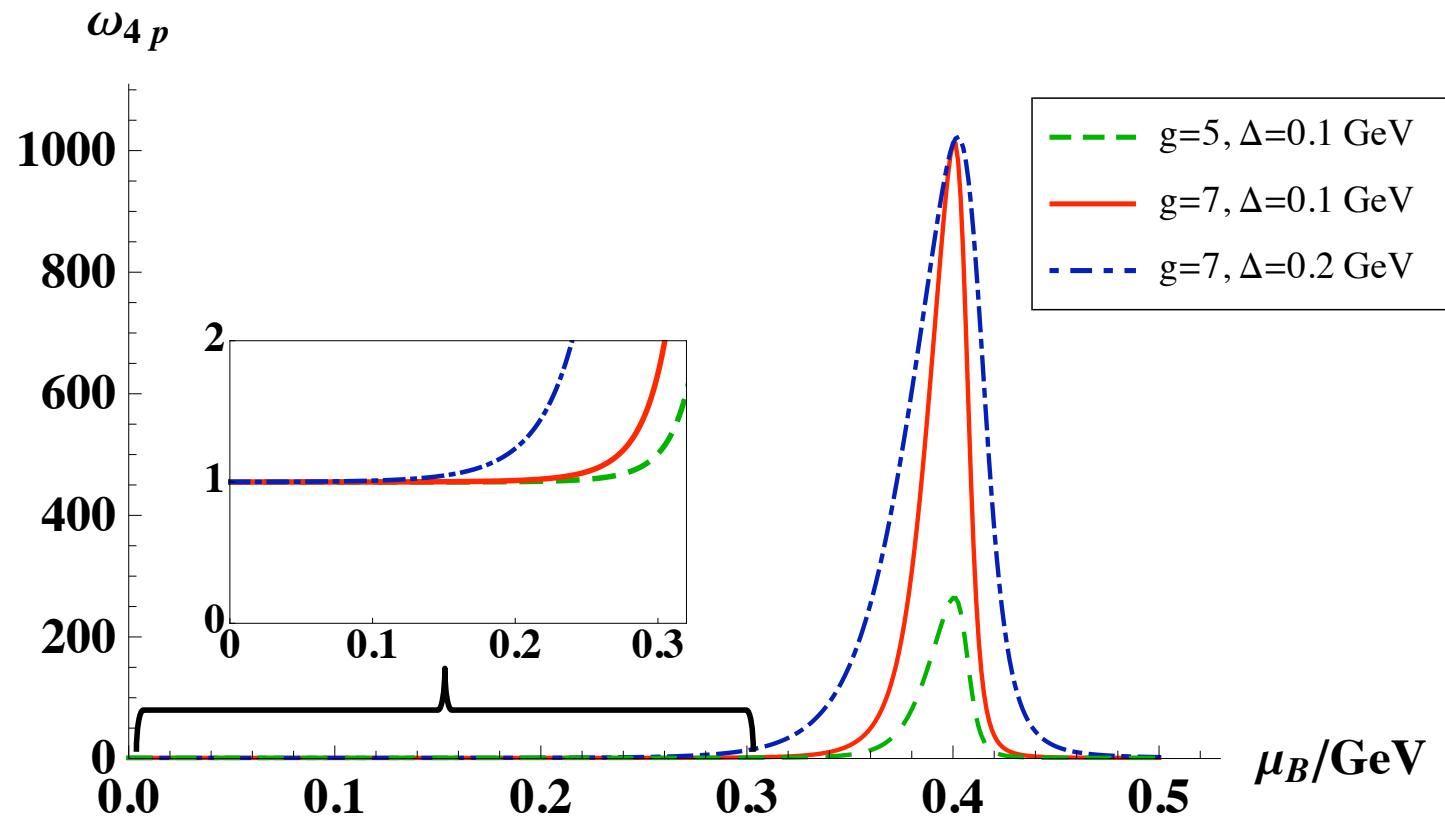
where  $\kappa \sigma^2 \equiv \frac{\kappa_4}{\kappa_2}$

(STAR Collaboration 2010)



Massachusetts  
Institute of  
Technology

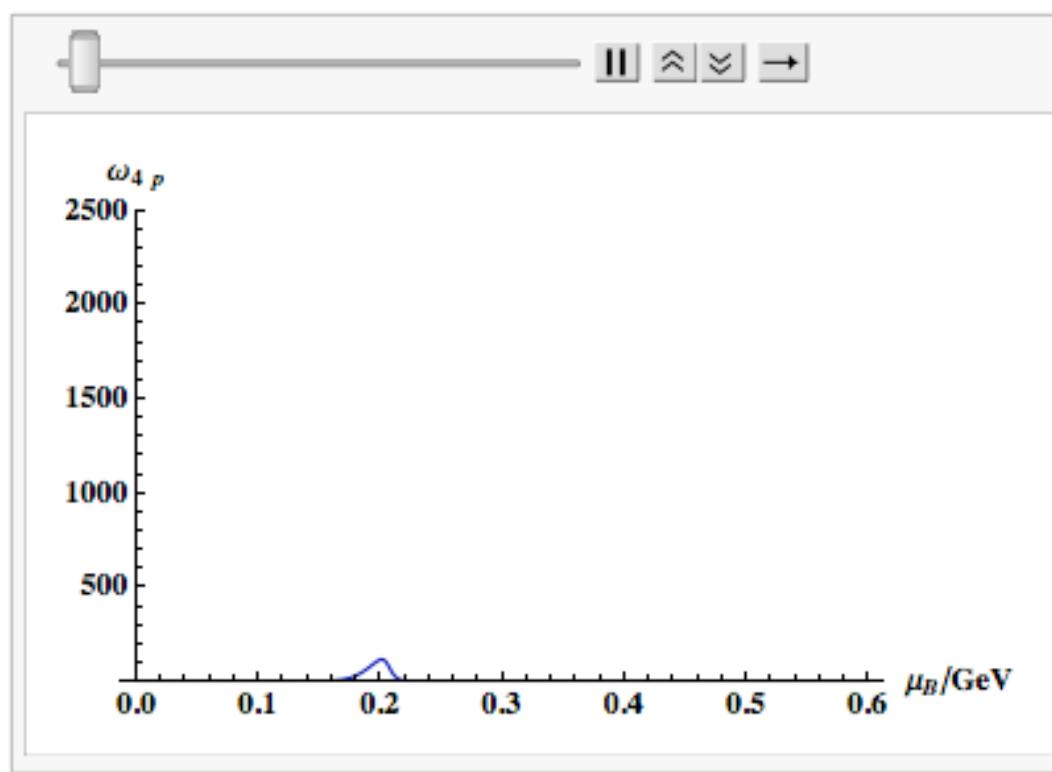
# Critical contribution to proton $\omega_4$



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# Multiplicity cumulants – movie

Changing the critical  $\mu_B$  – the location of the CP:



# Outline

- Introduction
- Critical Contribution to Particle Multiplicity Fluctuations
- Ratios of Fluctuation Observables
- Summary

# Uncertainties of parameters

- Cumulants depend on 5 non-universal parameters:

$$\kappa_{2\pi} \sim VT^{-1}G^2\xi^2N_\pi^2,$$

$$\kappa_{3\pi} \sim VT^{-3/2}G^3\tilde{\lambda}_3\xi^{9/2}N_\pi^3,$$

$$\kappa_{4\pi} \sim VT^{-2}G^4(2\tilde{\lambda}_3^2 - \tilde{\lambda}_4)\xi^7N_\pi^4$$



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- $G, g, \xi, \tilde{\lambda}_3, \tilde{\lambda}_4$  have large uncertainties
  - hard to predict the critical contribution to cumulants
- By taking ratios of cumulants can cancel some parameter dependence
  - minimize observable uncertainties

# Ratios of multiplicity cumulants

ratio	$V$	$n_p(\mu_B)$	$g$	$G$	$\tilde{\lambda}_3$	$\tilde{\lambda}'_4$	$\xi$
$N_\pi$	1	-	-	-	-	-	-
$N_p$	1	1	-	-	-	-	-
$\kappa_{ipj\pi}$	1	$i$	$i$	$j$	$\delta_{r,3}$	$\delta_{r,4}$	$\frac{5}{2}r - 3$
$\omega_{ipj\pi}$	-	$i - \frac{i}{r}$	$i$	$j$	$\delta_{r,3}$	$\delta_{r,4}$	$\frac{5}{2}r - 3$
$\kappa_{ipj\pi} N_\pi^{i-1}/N_p^i$	-	-	$i$	$j$	$\delta_{r,3}$	$\delta_{r,4}$	$\frac{5}{2}r - 3$
$\kappa_{2p2\pi} N_\pi / \kappa_{4\pi} \kappa_{2p}$	-	-	-	-2	-	-	-2
$\kappa_{4p} N_\pi^2 / \kappa_{4\pi} \kappa_{2p}^2$	-	-	-	-4	-	-	-4
$\kappa_{2p2\pi} N_p^2 / \kappa_{4p} N_\pi^2$	-	-	-2	2	-	-	-
$\kappa_{3p1\pi} N_p / \kappa_{4p} N_\pi$	-	-	-1	1	-	-	-
$\kappa_{3p} N_p^{3/2} / \kappa_{2p}^{9/4} N_\pi^{1/4}$	-	-	-3/2	-	1	-	-
$\kappa_{2p} \kappa_{4p} / \kappa_{3p}^2$	-	-	-	-	-2	1	-
$\kappa_{3p} N_\pi^{3/2} / \kappa_{3\pi} \kappa_{2p}^{3/2}$	-	-	-	-	-	-	-
$\kappa_{4p} N_\pi^{2/3} / \kappa_{4\pi} \kappa_{2p}^2$	-	-	-	-	-	-	-
$\kappa_{2p2\pi}^2 / \kappa_{4\pi} \kappa_{4p}$	-	-	-	-	-	-	-
$\kappa_{2p1\pi}^{2/3} N_\pi^{1/3} / \kappa_{3\pi}^2$	-	-	-	-	-	-	-

Ratios taken after subtracting Poisson and defined  $\tilde{\lambda}'_4 \equiv 2\tilde{\lambda}_3^2 - \tilde{\lambda}_4$

$$r = i + j$$



# Ratios of multiplicity cumulants

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$N_\pi$	1	-	-	-	-	-	-
$N_p$	1	1	-	-	-	-	-
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$\omega_{ipj\pi}$	-	$i - \frac{i}{r}$	$i$	$j$	$\delta_{r,3}$	$\delta_{r,4}$	$\frac{5}{2}r - 3$
$\kappa_{ipj\pi} N_\pi^{i-1}/N_p^i$	-	-	$i$	$j$	$\delta_{r,3}$	$\delta_{r,4}$	$\frac{5}{2}r - 3$
$\kappa_{2p2\pi} N_\pi / \kappa_{4\pi} \kappa_{2p}$	-	-	-	-2	-	-	-2
$\kappa_{4p} N_\pi^2 / \kappa_{4\pi} \kappa_{2p}^2$	-	-	-	-4	-	-	-4
$\kappa_{2p2\pi} N_p^2 / \kappa_{4p} N_\pi^2$	-	-	-2	2	-	-	-
$\kappa_{3p1\pi} N_p / \kappa_{4p} N_\pi$	-	-	-1	1	-	-	-
$\kappa_{3p} N_p^{3/2} / \kappa_{2p}^{9/4} N_\pi^{1/4}$	-	-	-3/2	-	1	-	-
$\kappa_{2p} \kappa_{4p} / \kappa_{3p}^2$	-	-	-	-	-2	1	-
$\kappa_{3p} N_\pi^{3/2} / \kappa_{3\pi} \kappa_{2p}^{3/2}$	-	-	-	-	-	-	-
$\kappa_{4p} N_\pi^{2/3} / \kappa_{4\pi} \kappa_{2p}^2$	-	-	-	-	-	-	-
$\kappa_{2p2\pi}^2 / \kappa_{4\pi} \kappa_{4p}$	-	-	-	-	-	-	-
$\kappa_{2p1\pi}^{2/3} N_\pi^{1/3} / \kappa_{3\pi}^2$	-	-	-	-	-	-	-

No parameter dependence

Ratios taken after subtracting Poisson and defined  $\tilde{\lambda}'_4 \equiv 2\tilde{\lambda}_3^2 - \tilde{\lambda}_4$

$$r = i + j$$



# Parameter independent ratios

- Parameter and energy independent ratios:

$$r_1 = \frac{\text{skewness}_p}{\text{skewness}_\pi}, \quad r_2 = \frac{\text{kurtosis}_p}{\text{kurtosis}_\pi}, \quad r_3 = \frac{\kappa_{2p1\pi}}{\kappa_{3\pi}^{2/3} \kappa_{3\pi}^{1/3}}, \quad r_4 = \frac{\kappa_{2p2\pi}^2}{\kappa_{4\pi} \kappa_{4p}}$$

where  $\text{skewness} = \frac{\kappa_3}{\kappa_2^{3/2}}$ ,  $\text{kurtosis} = \frac{\kappa_4}{\kappa_2^2}$

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- All equal to 1 if CP contribution dominates
- Poisson contribution:  $r_1 = (N_\pi/N_p)^{1/2}$ ,  $r_2 = N_\pi/N_p$ ,  $r_3 = r_4 = 0$
- How to use these ratios:
  - If one sees peaks in the measured cumulants at some  $\mu_B$
  - Calculate these ratios around the peak
  - If equal to 1  $\rightarrow$  Parameter independent way of verifying that the fluctuations you see are due to the CP

# Constraining parameters

- If CP found, can constrain parameters by measuring cumulant ratios near the CP

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- Parameters appear in certain combinations in the cumulants → can only constraint 4 independent (but not unique) combinations
- For example, some choices are:
  1.  $G \xi$  using  $\kappa_{2p2\pi} N_\pi / \kappa_{4\pi} \kappa_{2p}$  or  $\kappa_{4p} N_\pi^2 / \kappa_{4\pi} \kappa_{2p}^2$ ,
  2.  $G/g$  using  $\kappa_{2p2\pi} N_p^2 / \kappa_{4p} N_\pi^2$  or  $\kappa_{3p1\pi} N_p / \kappa_{4p} N_\pi$ ,
  3.  $\tilde{\lambda}'_4 / \tilde{\lambda}_3^2$  using  $\kappa_{2p} \kappa_{4p} / \kappa_{3p}^2$ ,
  4.  $\tilde{\lambda}_3^2 / g^3$  using  $\kappa_{3p} N_p^{3/2} / \kappa_{2p}^{9/4} N_\pi^{1/4}$ .

# Outline

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# Summary

- We used particle multiplicity fluctuations as a probe to the location of the CP
- Higher cumulants of event-by-event distributions are more sensitive to critical fluctuations
- CP signature: Non-monotonic behavior, as a function of collision energy, of multiplicity cumulants
- Constructed cumulant ratios to identify the CP location with reduced parameter uncertainties
- If CP is found, showed how to use cumulant ratios to constraint the values of the non-universal parameters



# Thank you!